QUANTUM ENTANGLEMENT AND QUANTUM INFORMATION

By AZMİ ALİ ALTINTAŞ OKAN UNIVERSITY

• EPR (Einstein-Podolsky-Rosen) Paper 1934



- EPR paper offered 2 explanations
- 1. Faster than light interaction (Einstein rejects this)
- 2. Hidden Variable

Hidden Variable theory states that there is a variable which has not been included QM yet, it contains all information.

• Bell's Theorem shows that 2nd argument is wrong.

There are four Bell pairs

• Bell Pairs

$$\begin{split} |\Phi^{+}\rangle_{AB} &= \frac{1}{\sqrt{2}} (|0\rangle_{A} \otimes |0\rangle_{B} + |1\rangle_{A} \otimes |1\rangle_{B}) \\ |\Phi^{-}\rangle_{AB} &= \frac{1}{\sqrt{2}} (|0\rangle_{A} \otimes |0\rangle_{B} - |1\rangle_{A} \otimes |1\rangle_{B}) \\ |\Psi^{+}\rangle_{AB} &= \frac{1}{\sqrt{2}} (|0\rangle_{A} \otimes |1\rangle_{B} + |1\rangle_{A} \otimes |0\rangle_{B}) \\ |\Psi^{-}\rangle_{AB} &= \frac{1}{\sqrt{2}} (|0\rangle_{A} \otimes |1\rangle_{B} - |1\rangle_{A} \otimes |0\rangle_{B}) \end{split}$$

• Other typical entangled states

1-GHZ (Greenberger-Horne-Zeilinger) state

$$|\mathrm{GHZ}\rangle = \frac{|0\rangle^{\otimes M} + |1\rangle^{\otimes M}}{\sqrt{2}}.$$

2-W state

$$|W\rangle = \frac{1}{\sqrt{n}}(|100...0\rangle + |010...0\rangle + ... + |00...01\rangle)$$

3-NOON state

$$|\psi_{\rm NOON}\rangle = \frac{|N\rangle_a|0\rangle_b + |0\rangle_a|N\rangle_b}{\sqrt{2}},$$

- Seperable states contains no entanglement
- All non-seperable states are entangled
- The entanglement of states does not increase under LOCC
- Entanglement does not change under local Unitary Operations
- There are maximally entangled states.



APPLICATION AREAS

- 1. Quantum Teleportation
- 2. Quantum Key Distribution
- 3. Quantum Computing

QUANTUM TELEPORTATION

• Quantum Teleportation is just a mechanism to "teleport" a qubit from one location to another one.

Alice and Bob are in different place and Alice wants to send an unkown qubit $|\psi\rangle_C = \alpha |0\rangle_C + \beta |1\rangle_C$. to Bob.

$$|\Phi^+\rangle_{AB} = \frac{1}{\sqrt{2}} (|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B)$$

From Bell states $|0\rangle \otimes |0\rangle = \frac{1}{\sqrt{2}}(|\Phi^+\rangle + |\Phi^-\rangle)$, etc. $|\Phi^+\rangle_{AB} \otimes |\psi\rangle_C = \frac{1}{2} \Big[|\Phi^+\rangle_{AC} \otimes (\alpha|0\rangle_B + \beta|1\rangle_B) + |\Phi^-\rangle_{AC} \otimes (\alpha|0\rangle_B - \beta|1\rangle_B) + |\Psi^+\rangle_{AC} \otimes (\beta|0\rangle_B + \alpha|1\rangle_B) + |\Psi^-\rangle_{AC} \otimes (\beta|0\rangle_B - \alpha|1\rangle_B) \Big].$

QUANTUM TELEPORTATION

Alice measures her state must be one of them

 $|\Phi^{+}\rangle_{AC}, |\Phi^{-}\rangle_{AC}, |\Psi^{+}\rangle_{AC}, |\Psi^{-}\rangle_{AC}$

She send her state to Bob via classical channel. Then Bob decides operations on his state.

- If $|\Phi^+\rangle_{AC}$, Bob does not do anyting.
- If $|\Phi^-\rangle_{AC}$ Bob does an operation $\sigma_3 = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix}$
- $|f|\Psi^+\rangle_{AC}$ Bob will do $\sigma_1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$

• For the remaining Bob does $-\sigma_3\sigma_1 = i\sigma_2 = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$.

QUANTUM KEY DISTRIBUTION

QKD guarantees secure communication. It is based on entanglement.(E91 Protocol by Artur Ekert 1991)

- 1. Entanglement establishes exact synchronization.
- 2. Any attemp at eavesdropping destroys correlation and it is detected.



Recipients can discern the presence of eavesdroppers because

QUANTUM COMPUTING

- Quantum computing studies on quantum computers that make direct use superposition and entanglement to perform operations on data.
- Instead of bits quantum bits (qubits) are used. $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$, Bloch sphere is a representation of qubit.



QUANTUM COMPUTING

- N qubits are the superposition of 2^N states
- Classical N bits just refers N states
- Quantum computers is used to simulate quantum many body systems

QUANTUM COMPUTING



512 qubit quantum computer developed by D-Wave.







FINE TUNNING-q deformation

- q-deformation is used to fit theoretical solutions to experimental results in atomic level.
- Deformed boson algebra

$$\begin{aligned} a_{q}a_{q}^{\dagger} - qa_{q}^{\dagger}a_{q} &= q^{-N}, \\ [a_{q}, N_{q}] &= a_{q}, \quad [a_{q}^{\dagger}, N_{q}] = -a_{q}^{\dagger} \\ [N]_{q} &= a_{q}^{\dagger}a_{q}, \quad a_{q}a_{q}^{\dagger} = [N+1]_{q}. \\ a_{q} &= \sqrt{\frac{q^{N}\psi_{1} - q^{-N}\psi_{2}}{N(q - q^{-1})}} a, \\ a_{q}^{\dagger} &= \sqrt{\frac{q^{N}\psi_{1} - q^{-N}\psi_{2}}{N(q - q^{-1})}} a^{\dagger}, \\ N_{q} &= N - (1/s)ln\psi_{2}. \end{aligned}$$

FINE TUNNING-q deformation

One can build deformed qubits and gates

Schwinger representation of angular momentum states

$$|jm\rangle = \frac{(a_1^{\dagger})^{j+m}(a_2^{\dagger})^{j-m}}{\sqrt{(j+m)!(j-m)!}}|\tilde{0_1}\tilde{0_2}\rangle.$$

$$\begin{split} |x\rangle_{q} &= (a_{1}^{\dagger})_{q}^{x} (a_{2}^{\dagger})_{q}^{(1-x)} |\tilde{0_{1}}\tilde{0_{2}}\rangle, \\ |1-x\rangle_{q} &= (a_{1}^{\dagger})_{q}^{1-x} (a_{2}^{\dagger})_{q}^{(x)} |\tilde{0_{1}}\tilde{0_{2}}\rangle. \end{split}$$

HNE IUNNING-q deformation

$$\begin{aligned} |x\rangle_{q} &= \left(\sqrt{\frac{q^{N_{1}}\psi_{1} - q^{-N_{1}}\psi_{2}}{N_{1}(q - q^{-1})}}a_{1}^{\dagger}\right)^{x}\left(\sqrt{\frac{q^{N_{2}}\psi_{3} - q^{-N_{2}}\psi_{4}}{N_{2}(q - q^{-1})}}a_{2}^{\dagger}\right)^{1-x}|\tilde{0_{1}}\tilde{0_{2}}\rangle, \\ |1-x\rangle_{q} &= \left(\sqrt{\frac{q^{N_{1}}\psi_{1} - q^{-N_{1}}\psi_{2}}{N_{1}(q - q^{-1})}}a_{1}^{\dagger}\right)^{1-x}\left(\sqrt{\frac{q^{N_{2}}\psi_{3} - q^{-N_{2}}\psi_{4}}{N_{2}(q - q^{-1})}}a_{2}^{\dagger}\right)^{x}|\tilde{0_{1}}\tilde{0_{2}}\rangle. \end{aligned}$$

HINE IUNNING-q deformation

$$Not_{q} = \sum_{x=0}^{1} |1 - x\rangle_{q} q\langle x|,$$

$$Had_{q} = -1^{(N_{1})} + \sum_{x=0}^{1} |1 - x\rangle_{q} q\langle x|$$

$$Swap_{q} = \sum_{x,y=0}^{1} |y x\rangle_{q} q\langle y x|,$$

$$Cnot_{q} = (1 - N_{1}) + \sum_{x,y=0}^{1} |x 1 - y\rangle_{q} q\langle y x|N_{1}$$

$$Fredkin_{q} = (1 - N_{1}) + \sum_{x,y,z=0}^{1} |x z y\rangle_{q} q\langle z y x|N_{1},$$

$$Toffoli_{q} = \sum_{x,y,z=0}^{1} |x y 1 - z\rangle_{q} q\langle z y x| [N_{1}M_{1} + (1 - N_{1})M_{1}]$$

$$+ \sum_{x,y,z=0}^{1} |x y 1 - z\rangle_{q} q\langle z y x| [(1 - N_{1})M_{1} + (1 - N_{1})(1 - M_{1})].$$

Constructing quantum logic gates using qdeformed harmonic oscillator algebras

ENTANGLEMENT MEASURES

- 1. CONCURRENCE
- 2. NEGATIVITY
- 3. RELATIVE ENTROPY OF ENTANGLEMENT
- 4. QUANTUM FISHER INFORMATION

CONCURRENCE

• Concurrence is an entanglement measure for a mixed state of two qubits

$$\mathcal{C}(\rho) \equiv \max(0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4)$$

 $\lambda_{1},...,$ λ_{4} are eigenvalues in descending order of a Hermitean matrix which is

$$R = \sqrt{\sqrt{\rho}\tilde{\rho}\sqrt{\rho}}$$

Where

$$\tilde{
ho} = (\sigma_y \otimes \sigma_y) \rho^* (\sigma_y \otimes \sigma_y)$$

If C is 0 the state is seperable, and maximally entangled for 1.

NEGATIVITY

• $N(\rho) = max\{0, -2\mu_{min}\}$

 $\mu_{\rm min}$ is the minimum eigenvalue of the partial transpose of the density matrix ρ

It can takes values between 0 and 1. If $N(\rho)$ is 0 the system is separable, and if it is 1 the system is maximally entangled as in concurrence.

RELATIVE ENTROPY OF ENTANGLEMENT

$E_R(\rho) = \min_{\sigma' \in D} S(\rho || \sigma') = S(\rho || \sigma)$

D is a set of separable states. S is quasidistance mesure

$S(\rho || \sigma') = \operatorname{tr} \left(\rho \lg \rho - \rho \lg \sigma'\right)$

 σ is a state on the boundary of separable states which is called closest separable state. If E is 0 the state is separable.

RELATIVE ENTROPY OF ENTANGLEMENT



- Quantum Fisher Information (QFI) is related with estimating a parameter θ by making measurement on the operator A.
- How much precise can a phase be estimated.
- Cramér-Rao bound gives a lower bound on the pressure $(\Delta \theta)^2 \ge \frac{1}{F_Q[\varrho, A]}$ estimation

A is any Hermitean operator

- Entanglement can increase the sensitivity of an interferometer.
- Example: Mach-Zehnder interferometer



• For non entangled states $(\Delta \theta)^2 \sim \frac{1}{N}$ called as shot noise limit

• Quantum entanglement makes it possible to reach

$$(\Delta\theta)^2 \sim \frac{1}{N^2}$$

which is the Heisenberg limit

Quantum Fisher Information is defined as

 $F(\theta) = tr\rho(\theta)LL^{\dagger}$

Quantum score has properties

$$\frac{\partial}{\partial \theta} \rho(\theta) = \rho(\theta)L \text{ right}$$
$$\frac{\partial}{\partial \theta} \rho(\theta) = L \rho(\theta) \text{ left}$$
$$\frac{\partial}{\partial \theta} \rho(\theta) = \frac{1}{2} \left(\rho(\theta)L + L \rho(\theta) \right) \text{ symmetric}$$

From the general expression of QFI one can derive many formulas for QFI by using definitions of quantum score.

In our works we use

$$F(\rho, J_{\overrightarrow{n}}) = \sum_{i \neq j} \frac{2(p_i - p_j)^2}{p_i + p_j} |\langle i | J_{\overrightarrow{n}} | j \rangle|^2 = \overrightarrow{n} C \overrightarrow{n}^T.$$

$$C_{kl} = \sum_{i \neq j} \frac{(p_i - p_j)^2}{p_i + p_j} [\langle i | J_k | j \rangle \langle j | J_l | i \rangle + \langle i | J_l | j \rangle \langle j | J_k | i \rangle].$$

here p_i and |j> are the eigenvalue and eigenvector of $\vec{F}_{max} = \frac{\hat{F}_{max}}{N} = \frac{\lambda_{max}}{N}$

Quantum Fisher Information

• QFI for seperable states

$$\overline{F}_{max} \leq 1$$

For general states

$\overline{F}_{max} \leq N$

For maximally entangled states the upper limit is saturated.

QFI is not an entanglement measure!!

J. Opt. B: Quantum Semiclass. Opt. 6 (2004) 542-548

PII: S1464-4266(04)86422-X

A comparative study of relative entropy of entanglement, concurrence and negativity

Adam Miranowicz and Andrzej Grudka



SCIENTIFIC REPORTS

www.nature.com/**scientificreports**

OPEN

SUBJECT AREAS: COMPUTER SCIENCE QUANTUM INFORMATION QUBITS

> Received 4 April 2014

Accepted 4 June 2014 Published

24 June 2014

Systems Volkan Erol^{1,2}, Fatih Ozaydin³ & Azmi Ali Altintas⁴

Analysis of Entanglement Measures and

LOCC Maximized Quantum Fisher

Information of General Two Qubit

¹Institute of Science, Okan University, Istanbul, Turkey, ²Provus A MasterCard Company R&D Center, Istanbul, Turkey, ³Department of Information Technologies, Isik University, Istanbul, Turkey, ⁴Department of Electrical Engineering, Okan University, Istanbul, Turkey.

Since QFI changes under unitary operations, one can find a unitary transformation in which QFI gives the same results with entanglement measures.

• An Euler rotation is defined on each qubit

$$U_{Rot}(\alpha,\beta,\gamma) = U_x(\alpha)U_z(\beta)U_x(\gamma)$$

Where the unitary rotations

$$U_j(\alpha) = \exp\left(-i\alpha \frac{\sigma_j}{2}\right), j \in \{x,z\}$$



- Maximal QFI is just the maximum eigenvalue of QFI matrix.
- Maximized QFI by LOCC behaves similar to the entanglement measures.
- If the rotation angle is chosen as smaller than 120 degree (which is used in this study) we can get better result.

Papers

- 1. V. Erol, F. Özaydın, A. A. Altıntaş, "Scientific Reports", 4, June 2014.
- 2. F. Özaydın, S. Bugu, C. Yeşilyurt, A. A. Altıntaş, M. Tame, Ş.K. Özdemir, "Physical Review A", 89, April-2014.
- 3. A.A Altintas, F. Ozaydin, C. Yesilyurt, S. Bugu and M. Arık "QIP", 13, 2014.
- 4. F. Ozaydin, A. A. Altintas, S. Bugu, C. Yesilyurt, "IJTP" 52-9, 2013.
- 5. F. Ozaydin, A. A. Altintas, S. Bugu, C. Yesilyurt and M. Arik "IJTP" 53-5, 2014.

Papers

- 6. S. Bugu, C. Yesilyurt, F. Diker, A. A. Altintas, F. Ozaydin "Acta Physica Polonica A" 127,2015.
- 7. F. Ozaydin, A. A. Altintas, S. Bugu, C. Yesilyurt andV. Erol "Acta Physica Polonica A" 127,2015.
- 8. F. Ozaydin, A. A. Altintas, S. Bugu, C. Yesilyurt "Acta Physica Polonica A" 124,2014.
- 9. F. Ozaydin, A. A. Altintas, "Scientific Reports", 2015.
- 10. A. A. Altintas, "Annals of Physics", 2016.

THANKS